

***P22.7** For each electron, $|q|vB \sin 90.0^\circ = \frac{mv^2}{r}$ and $v = \frac{eBr}{m}$.

The electrons have no internal structure to absorb energy, so the collision must be perfectly elastic:

$$K = \frac{1}{2}mv_{1i}^2 + 0 = \frac{1}{2}mv_{1f}^2 + \frac{1}{2}mv_{2f}^2$$

$$K = \frac{1}{2}m \left(\frac{e^2 B^2 R_1^2}{m^2} \right) + \frac{1}{2}m \left(\frac{e^2 B^2 R_2^2}{m^2} \right) = \frac{e^2 B^2}{2m} (R_1^2 + R_2^2)$$

$$K = \frac{e(1.60 \times 10^{-19} \text{ C})(0.0440 \text{ N} \cdot \text{s/C} \cdot \text{m})^2}{2(9.11 \times 10^{-31} \text{ kg})} [(0.0100 \text{ m})^2 + (0.0240 \text{ m})^2] = \boxed{115 \text{ keV}}$$

P22.10 $F_B = F_e$

so $qvB = qE$

where $v = \sqrt{\frac{2K}{m}}$ and K is kinetic energy of the electron.

$$E = vB = \sqrt{\frac{2K}{m}} B = \sqrt{\frac{2(750)(1.60 \times 10^{-19})}{9.11 \times 10^{-31}}} (0.0150) = \boxed{244 \text{ kV/m}}$$

P22.13 $\theta = \tan^{-1} \left(\frac{25.0}{10.0} \right) = 68.2^\circ$ and $R = \frac{1.00 \text{ cm}}{\sin 68.2^\circ} = 1.08 \text{ cm}$.

Ignoring relativistic correction, the kinetic energy of the electrons is

$$\frac{1}{2}mv^2 = q\Delta V \quad \text{so} \quad v = \sqrt{\frac{2q\Delta V}{m}} = 1.33 \times 10^8 \text{ m/s}.$$

From Newton's second law $\frac{mv^2}{R} = qvB$, we find the magnetic field

$$B = \frac{mv}{|q|R} = \frac{(9.11 \times 10^{-31} \text{ kg})(1.33 \times 10^8 \text{ m/s})}{(1.60 \times 10^{-19} \text{ C})(1.08 \times 10^{-2} \text{ m})} = \boxed{70.1 \text{ mT}}.$$

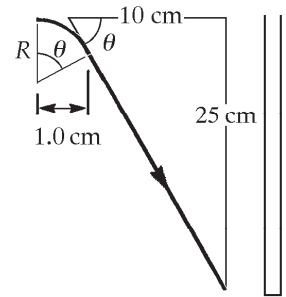


FIG. P22.13

P22.14 (a) If the charge carriers are negative, to carry current in the x direction they move with drift velocity v_d in the $-x$ direction. The magnetic force $q\vec{v} \times \vec{B}$ is in the $-(\hat{i}) \times \hat{j} = \hat{k}$ direction, so the negative charges are deflected to the top of the ribbon, point c . Some accumulate there to make V_c negative with respect to V_a , until ...

- (b) ... an upward electric field $\vec{E} = \frac{|V_c - V_a|}{d} \hat{\mathbf{k}}$ exerts a downward force on the other charge carriers to let them drift in equilibrium according to

$$\Sigma F_z = 0: \quad |q| \frac{|V_c - V_a|}{d} (-\hat{\mathbf{k}}) + |q| |\vec{v}_d| |\vec{B}| \hat{\mathbf{k}} = 0$$

$$v_d = \frac{|\Delta V_H|}{dB}$$

Since the measured current is $I = n |q| v_d t d$

we have $I = n |q| \frac{|\Delta V_H|}{dB} t d$

$$n = \frac{IB}{|q| |\Delta V_H| t}$$

Since we have shown that ΔV_H is negative if q is negative, this expression simplifies to

$$n = \frac{IB}{q \Delta V_H t}$$

P22.15 $\vec{F}_B = I \vec{\ell} \times \vec{B} = (2.40 \text{ A})(0.750 \text{ m}) \hat{\mathbf{i}} \times (1.60 \text{ T}) \hat{\mathbf{k}} = \boxed{(-2.88 \hat{\mathbf{j}}) \text{ N}}$

- P22.17** The magnetic force on each bit of ring is $I d\vec{s} \times \vec{B} = IdsB$ radially inward and upward, at angle θ above the radial line. The radially inward components tend to squeeze the ring but all cancel out as forces. The upward components $IdsB \sin \theta$ all add to $\boxed{I2\pi rB \sin \theta \text{ up}}$.

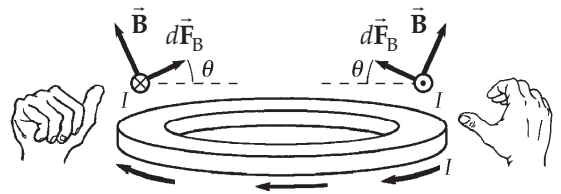
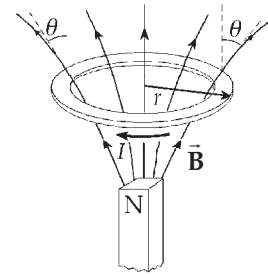


FIG. P22.17

P22.19 (a) $2\pi r = 2.00 \text{ m}$

so $r = 0.318 \text{ m}$

$$\mu = IA = (17.0 \times 10^{-3} \text{ A}) [\pi (0.318)^2 \text{ m}^2] = \boxed{5.41 \text{ mA} \cdot \text{m}^2}$$

$$(b) \quad \vec{\tau} = \vec{\mu} \times \vec{B}$$

$$\text{so} \quad \tau = (5.41 \times 10^{-3} \text{ A} \cdot \text{m}^2)(0.800 \text{ T}) = \boxed{4.33 \text{ mN} \cdot \text{m}}$$

$$\text{P22.23} \quad B = \frac{\mu_0 I}{2R} = \frac{\mu_0 q(v/2\pi R)}{2R} = \boxed{12.5 \text{ T}}$$

P22.27 We can think of the total magnetic field as the superposition of the field due to the long straight wire (having magnitude $\frac{\mu_0 I}{2\pi R}$ and directed into the page) and the field due to

the circular loop (having magnitude $\frac{\mu_0 I}{2R}$ and directed into the page). The resultant

$$\text{magnetic field is: } \boxed{\vec{B} = \left(1 + \frac{1}{\pi}\right) \frac{\mu_0 I}{2R} \text{ (directed into the page)}}$$

***P22.29** Wire 1 creates at the origin magnetic field

$$\vec{B}_1 = \frac{\mu_0 I}{2\pi r} \text{ right hand rule} = \frac{\mu_0 I_1}{2\pi a} \hat{j} = \frac{\mu_0 I_1}{2\pi a} \hat{j}$$

(a) If the total field at the origin is $\frac{2\mu_0 I_1}{2\pi a} \hat{j} = \frac{\mu_0 I_1}{\pi a} \hat{j} + \vec{B}_2$ then the second wire must

$$\text{create field according to } \vec{B}_2 = \frac{\mu_0 I_1}{2\pi a} \hat{j} = \frac{\mu_0 I_2}{2\pi(2a)} \hat{j}$$

$$\text{Then } I_2 = \boxed{2I_1 \text{ out of the paper}}$$

(b) The other possibility is $\vec{B}_1 + \vec{B}_2 = \frac{2\mu_0 I_1}{2\pi a} (-\hat{j}) = \frac{\mu_0 I_1}{\pi a} (-\hat{j}) + \vec{B}_2$. Then

$$\vec{B}_2 = \frac{3\mu_0 I_1}{2\pi a} (-\hat{j}) = \frac{\mu_0 I_2}{2\pi(2a)} (-\hat{j}) \quad I_2 = \boxed{6I_1 \text{ into the paper}}$$

P22.31 We use the Biot-Savart law. For bits of wire along the straight-line sections, $d\vec{s}$ is at 0° or 180° to \hat{r} , so $d\vec{s} \times \hat{r} = 0$. Thus, only the curved section of wire contributes to \vec{B} at P . Hence, $d\vec{s}$ is tangent to the arc and \hat{r} is radially inward; so $d\vec{s} \times \hat{r} = |ds| \ell \sin 90^\circ = |ds| \otimes$. All points along the curve are the same distance $r = 0.600 \text{ m}$ from the field point, so

$$B = \int_{\text{all current}} |d\vec{B}| = \int \frac{\mu_0 I}{4\pi} \frac{|d\vec{s} \times \hat{r}|}{r^2} = \frac{\mu_0 I}{4\pi r^2} \int |ds| = \frac{\mu_0 I}{4\pi r^2} s$$

where s is the arc length of the curved wire,

$$s = r\theta = (0.600 \text{ m})(30.0^\circ) \left(\frac{2\pi}{360^\circ}\right) = 0.314 \text{ m}$$

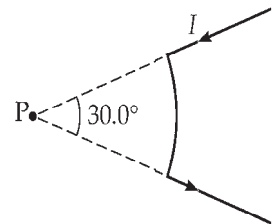


FIG. P22.31

Then, $B = (10^{-7} \text{ T} \cdot \text{m/A}) \frac{(3.00 \text{ A})}{(0.600 \text{ m})^2} (0.314 \text{ m})$

$B = \boxed{261 \text{ nT into the page}}$

P22.34 Let both wires carry current in the x direction, the first at $y = 0$ and the second at $y = 10.0 \text{ cm}$.

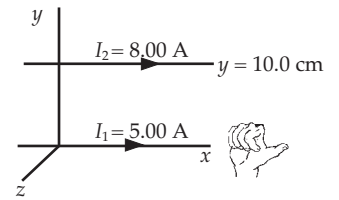


FIG. P22.34(a)

(a) $\vec{B} = \frac{\mu_0 I}{2\pi r} \hat{k} = \frac{(4\pi \times 10^{-7} \text{ T} \cdot \text{m/A})(5.00 \text{ A})}{2\pi(0.100 \text{ m})} \hat{k}$

$\vec{B} = \boxed{1.00 \times 10^{-5} \text{ T out of the page}}$

(b) $\vec{F}_B = I_2 \vec{\ell} \times \vec{B} = (8.00 \text{ A})[(1.00 \text{ m})\hat{i} \times (1.00 \times 10^{-5} \text{ T})\hat{k}] = (8.00 \times 10^{-5} \text{ N})(-\hat{j})$

$\vec{F}_B = \boxed{8.00 \times 10^{-5} \text{ N toward the first wire}}$

(c) $\vec{B} = \frac{\mu_0 I}{2\pi r} (-\hat{k}) = \frac{(4\pi \times 10^{-7} \text{ T} \cdot \text{m/A})(8.00 \text{ A})}{2\pi(0.100 \text{ m})} (-\hat{k}) = (1.60 \times 10^{-5} \text{ T})(-\hat{k})$

$\vec{B} = \boxed{1.60 \times 10^{-5} \text{ T into the page}}$

continued on next page

(d) $\vec{F}_B = I_1 \vec{\ell} \times \vec{B} = (5.00 \text{ A})[(1.00 \text{ m})\hat{i} \times (1.60 \times 10^{-5} \text{ T})(-\hat{k})] = (8.00 \times 10^{-5} \text{ N})(+\hat{j})$

$\vec{F}_B = \boxed{8.00 \times 10^{-5} \text{ N towards the second wire}}$

P22.35 By symmetry, we note that the magnetic forces on the top and bottom segments of the rectangle cancel. The net force on the vertical segments of the rectangle is (using Equation 22.27)

$$\vec{F} = \vec{F}_1 + \vec{F}_2 = \frac{\mu_0 I_1 I_2 \ell}{2\pi} \left(\frac{1}{c+a} - \frac{1}{c} \right) \hat{i} = \frac{\mu_0 I_1 I_2 \ell}{2\pi} \left(\frac{-a}{c(c+a)} \right) \hat{i}$$

$$\vec{F} = \frac{(4\pi \times 10^{-7} \text{ N/A}^2)(5.00 \text{ A})(10.0 \text{ A})(0.450 \text{ m})}{2\pi} \left(\frac{-0.150 \text{ m}}{(0.100 \text{ m})(0.250 \text{ m})} \right) \hat{i}$$

$$\vec{F} = (-2.70 \times 10^{-5} \hat{i}) \text{ N}$$

or $\vec{F} = \boxed{2.70 \times 10^{-5} \text{ N toward the left}}$

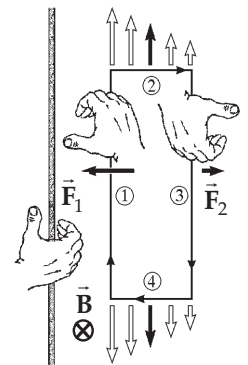


FIG. P22.35